

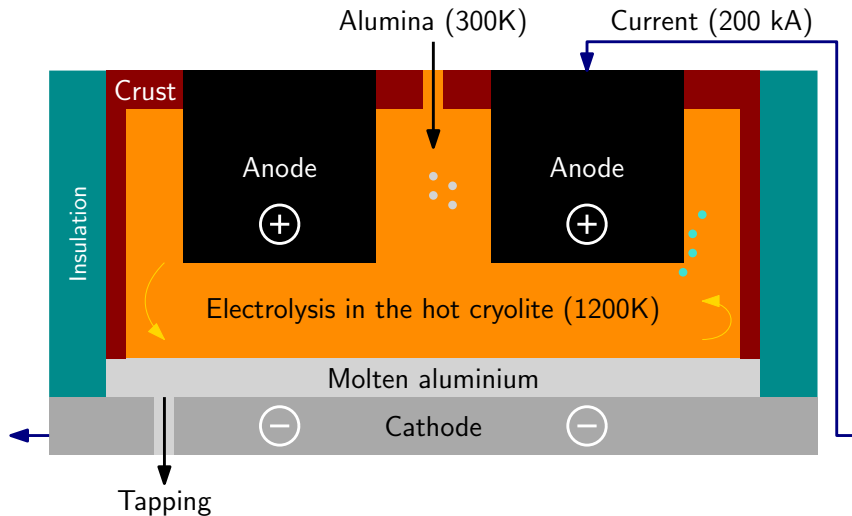
# Mathematical modelling of alumina feeding

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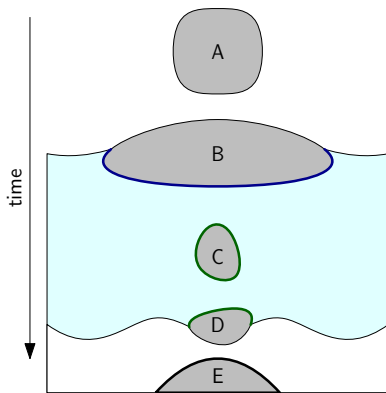
Chris Breward, James Oliver and Andreas Münch (Oxford)  
Svenn Halvorsen and Ellen Nordgård-Hansen (Norce)  
Eirik Manger (Norsk Hydro)

March 14, 2019

## Hall-Héroult cell

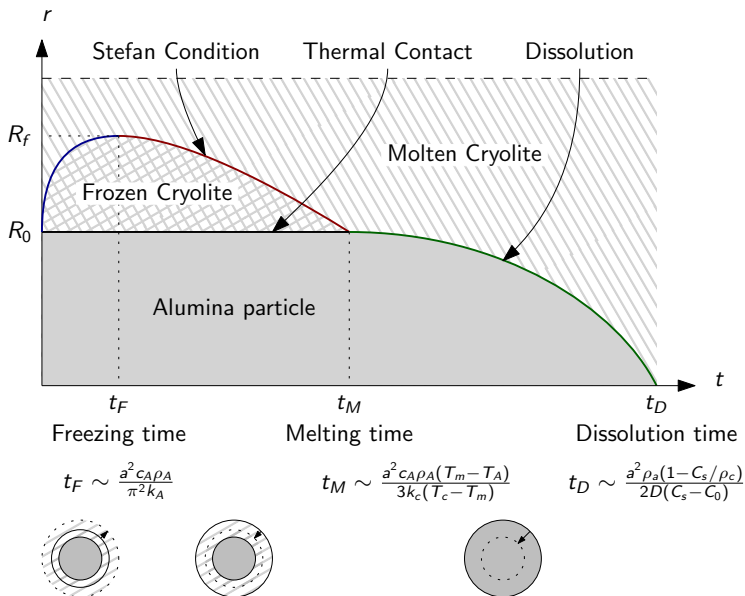


## Alumina feeding processes

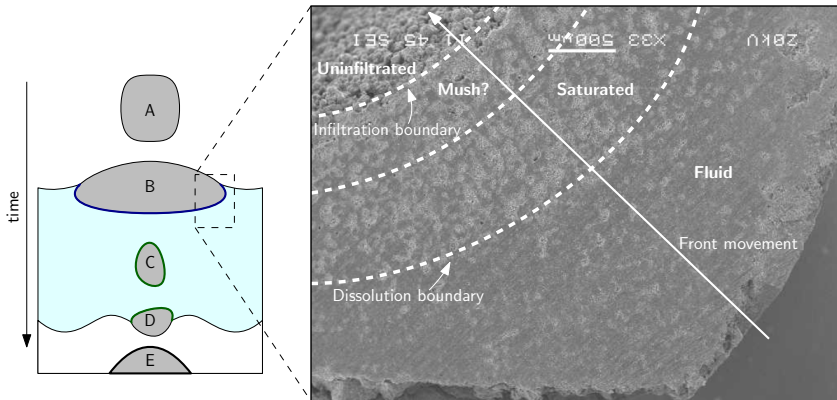


How does the molten cryolite infiltrate and dissolve a porous alumina structure?

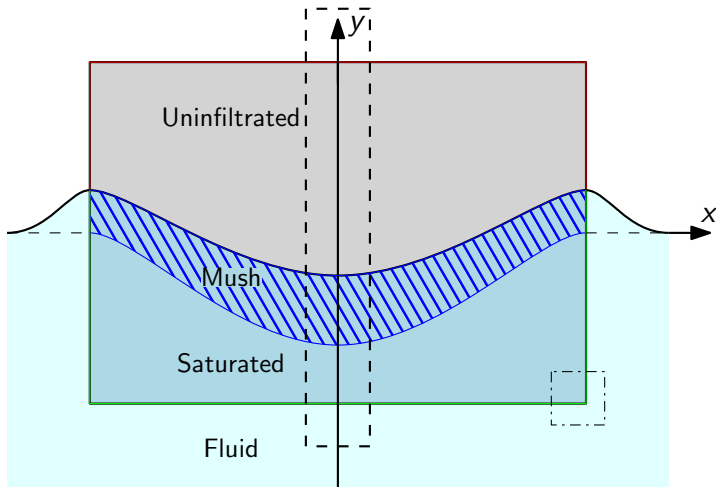
## Previous Analysis: Solid alumina particle



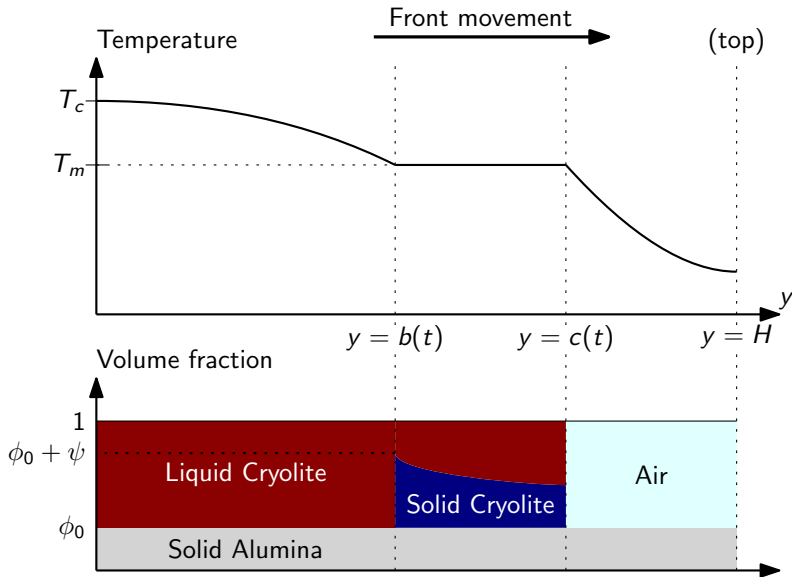
# Alumina feeding: raft problem



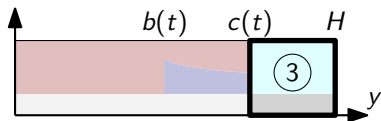
## Alumina feeding: raft problem



## 1D Infiltration problem



## 1D Infiltration problem: uninfiltrated region



For  $c(t) < y < H$ :

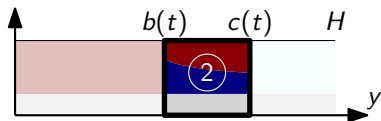
$$\frac{\partial T_3}{\partial t} = \frac{k^*}{\rho_A c_A \phi} \frac{\partial^2 T_3}{\partial y^2}$$

$$\text{At } y = c(t): \quad T_3 = T_m, \quad \rho_c \psi \dot{c} = -k_3 \frac{\partial T_3}{\partial y}$$

$$\text{At } y = H: \quad k_3 \frac{\partial T_3}{\partial y} = h(T - T_e)$$



## 1D Infiltration problem: mushy region



For  $b(t) < y < c(t)$ :

$$\frac{\partial \psi}{\partial t} = 0$$

$$\frac{\partial}{\partial y} ((1 - \phi - \psi)u_2) = 0$$

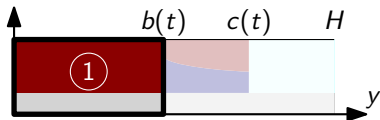
$$\frac{-K(\psi + \phi)}{\mu(1 - \phi - \psi)} \left( \frac{\partial p_2}{\partial y} + \rho_c g \right) = u_2$$

$$\text{At } y = c(t) : \quad p_2 = p_a - \frac{\sigma_c}{r_{pore}}, \quad u_2 = \dot{c} \left( 1 + \frac{\psi \rho}{1 - \phi - \psi} \right)$$

$$\text{At } y = b(t) : \quad [p_i] = 0, \quad u_1 - \frac{1 - \phi - \psi}{1 - \phi} u_2 = \psi \dot{b} \frac{1 - \rho}{1 - \phi}$$

## 1D Infiltration problem: infiltrated region

For  $0 < y < b(t)$ :



$$\frac{\partial T_1}{\partial t} + \frac{\partial}{\partial y} \left( \frac{1 - \phi}{a} u_1 T_1 \right) = \frac{k_1}{a \rho_A c_A} \frac{\partial^2 T}{\partial y^2}$$

$$\frac{\partial}{\partial y} ((1 - \phi) u_1) = 0$$

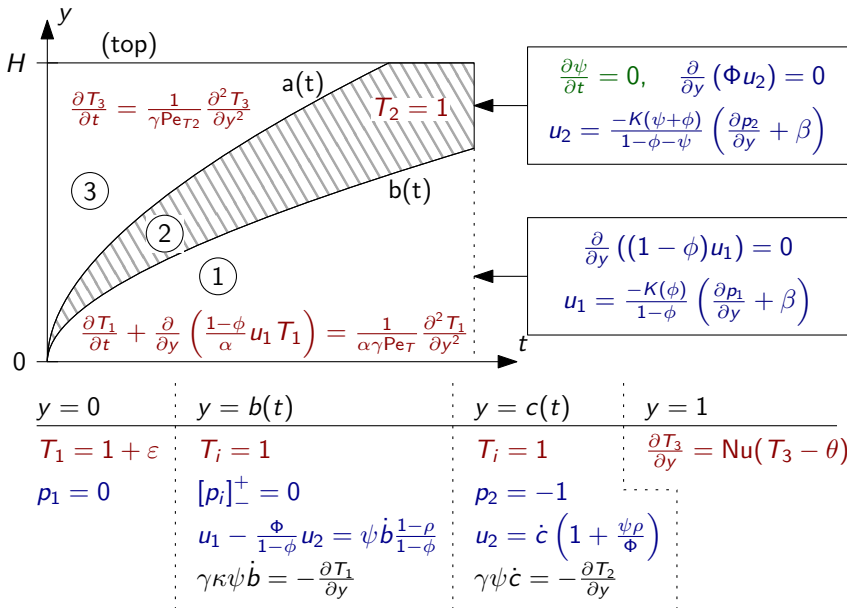
$$\frac{-K(\phi)}{\mu(1 - \psi)} \left( \frac{\partial p_1}{\partial y} + \rho_c g \right) = u_1$$

$$\text{At } y = b(t): \quad T_1 = T_m, \quad \rho_c \psi \dot{b} = -k_1 \frac{\partial T}{\partial y},$$

$$[p_i] = 0, \quad u_1 - \frac{1 - \phi - \psi}{1 - \phi} u_2 = \psi \dot{b} \frac{1 - \rho}{1 - \phi},$$

$$\text{At } y = 0: \quad T_1 = T_h, \quad p_1 = \rho g H$$

## Full dimensionless problem



## Dimensionless parameters

Meaning	Symbol	Equation	Value
Infiltrating force	$\gamma$	$KP_T/(\mu_c H^2)$	40–400
Newton cooling	Nu	$hH/k^*$	3
Conductivity ratio	$\kappa$	$k^*/k_m$	0.2
Thermal diffusivity	$Pe_T$	$H^2 \rho_c c_c / (k_m \tau)$	1.3
Thermal diffusivity	$Pe_{T2}$	$\frac{H^2 \rho_{ACA} \phi}{\tau k_A (k_{air}/k_A)^{1-\phi}}$	1.77
Particle coldness	$\theta$	$T_A/T_m$	0.246
Overheat	$\varepsilon_T$	$T_c/T_m - 1$	0.015
Density ratio	$\rho$	$\rho_f/\rho_c$	1.00966
Importance of gravity	$\beta$	$\rho_c g H / P$	0.01
Energy stored in phases	$\alpha$	$\phi \frac{\rho_{ACA}}{\rho_c c_c} + (1 - \phi)$	0.81

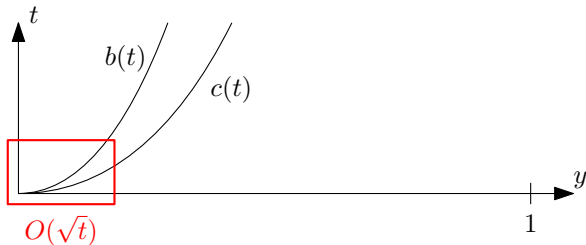
# Small time asymptotics

- Expect  $b, c = O(\sqrt{t})$  as  $t \rightarrow 0^+$ , so need small  $t$  asymptotics to initiate numerical solution.
- Seek similarity solution given

$$T_i \sim f_i(\eta), \quad u_i \sim \frac{g_i(\eta)}{\sqrt{t}}, \quad p_i \sim h_i(\eta), \quad \psi \sim \psi_c,$$

$$b \sim 2\lambda_b\sqrt{t}, \quad c \sim 2\lambda_c\sqrt{t}$$

with  $\eta = y/2\sqrt{t} = O(1)$  as  $t \rightarrow 0^+$  and  $\lambda_b, \lambda_c, \psi_c$  to be determined.



# Transcendental system for $\lambda_b, \lambda_c, \psi_c$

$$K_2(\psi_c) = -2(\lambda_b - \lambda_c)g_2(1 - \phi - \psi_c) + \lambda_b g_1(1 - \phi) \frac{K_2(\psi_c)}{K_1}$$

$$\gamma 2\kappa \psi_c \lambda_b = \frac{2\sqrt{\text{Pe}_{T1}} \epsilon e^{-\text{Pe}_{T1}(ag_1 + \lambda_b)^2}}{\sqrt{\pi} (\text{erf}(\sqrt{\text{Pe}_{T1}}(ag_1 + \lambda_b)) - \text{erf}(ag_1\sqrt{\text{Pe}_{T1}}))}$$

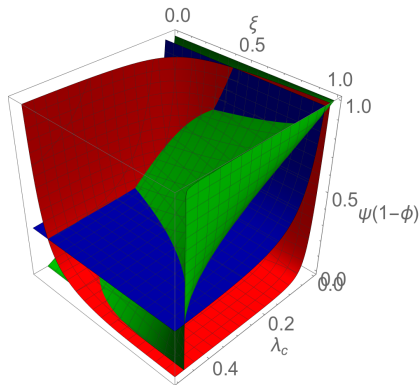
$$2\gamma \psi_c \lambda_c = -\frac{2(\theta - 1)\sqrt{\text{Pe}_{T2}} e^{-\text{Pe}_{T2}\lambda_c^2}}{\sqrt{\pi} \text{erfc}(\lambda_c\sqrt{\text{Pe}_{T2}})}$$

where  $g_1$  and  $g_2$  constants given by

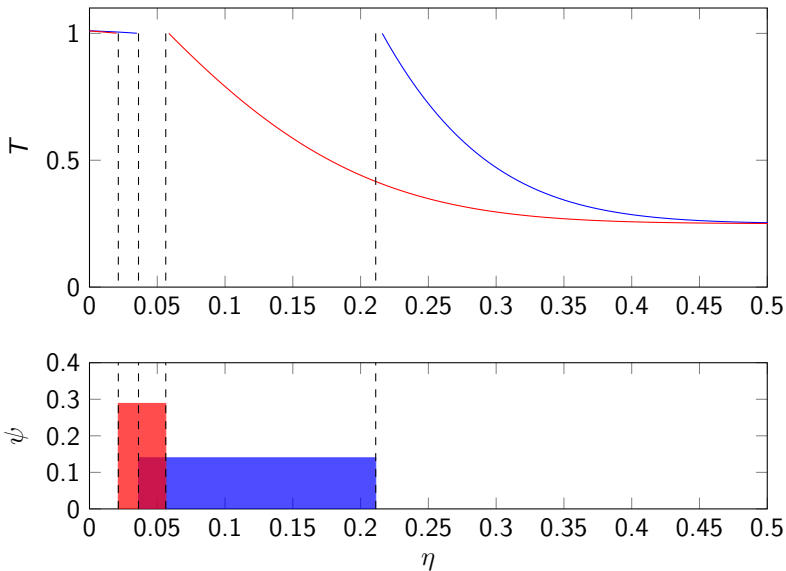
$$g_1 = \lambda_c + \frac{(\lambda_c - \lambda_b)(\rho - 1)\psi_c}{1 - \phi}$$

$$g_2 = \lambda_c \left( 1 + \frac{\psi_c \rho}{1 - \phi - \psi_c} \right)$$

## Constraint surfaces

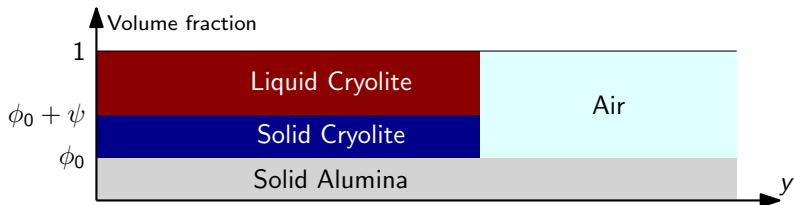
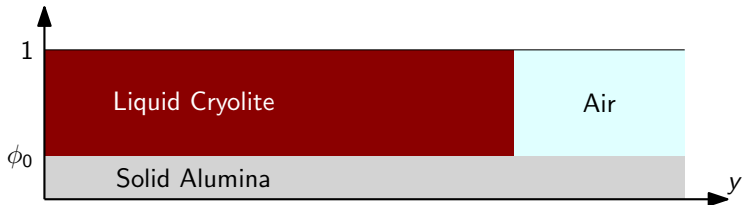


Two physical solutions with  $\gamma = 200$  and  $\text{Pe}_{T2} = 1/7$ .

Similarity solution for  $\gamma = 200$ ,  $Pe_{T2} = 0.14$ 

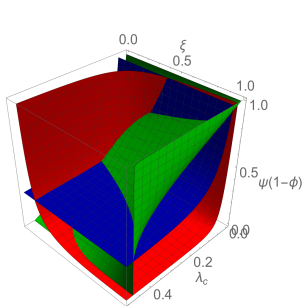


## Sanity check: sublimits

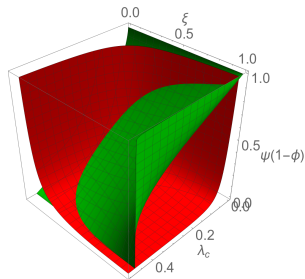
No heating ( $\varepsilon = 0, \lambda_b = 0$ )No cooling and heating ( $\varepsilon, \theta = 1$ )

# Comparison with numerical results

## Nonexistence of Solutions



No physical solutions with  $\gamma = 75$  and  $Pe_{T2} = 0.14$  due to no intersection in the physically admissible region.



No physical solutions with  $\gamma = 75$  and  $Pe_{T2} = 1.7$  due to the Stefan condition (blue).  $\psi > 1 - \phi$ , meaning freezes more than the available fluid.

# Conclusions and Future work

## Conclusions:

- Developed a simple model for the infiltration of molten cryolite into a cold porous alumina structure (with phase change).
- Similarity solution at small times yields nonexistence and nonuniqueness.
- Preliminary numerical simulation suggest one solution is admissible while the other may blow up.

## Future work:

- Characterise regions for similarity solution
- Investigate small time behaviour numerically: Alternatives to similarity solution?
- Are we missing physics e.g.
  - capillary pressure dependence on porosity
  - dropping local thermodynamical equilibrium
  - including composition effects